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B.Sc. Part II (Hons.) 4th Paper

DYNAMICS

Q. A particle moves along a circle $r = 2a \cos \theta$ in such a way that its acceleration towards the origin is always zero. Prove that $\frac{d\omega}{dt} = -2\omega^2 \cot \theta$, where $\omega = \frac{d\theta}{dt}$.

Soln Given that acceleration towards origin = 0
ie. radial acceleration = 0.

$$\Rightarrow \frac{d^2 r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 = 0 \quad \text{But } \omega = \frac{d\theta}{dt}$$

$$\Rightarrow \frac{d^2 r}{dt^2} - r \omega^2 = 0 \Rightarrow \frac{d^2 r}{dt^2} = r \omega^2 \quad \text{--- (1)}$$

Also, given $r = 2a \cos \theta$

$$\Rightarrow \frac{dr}{dt} = -2a \sin \theta \frac{d\theta}{dt}$$

$$\Rightarrow \frac{d^2 r}{dt^2} = -2a \cos \theta \left(\frac{d\theta}{dt} \right)^2 - 2a \sin \theta \frac{d^2 \theta}{dt^2} \quad \text{--- (2)}$$

$$\text{But } \frac{d\omega}{dt} = \frac{d}{dt} \left(\frac{d\theta}{dt} \right) = \frac{d^2 \theta}{dt^2}$$

$$\text{So, (2)} \Rightarrow \frac{d^2 r}{dt^2} = -2a \cos \theta \left(\frac{d\theta}{dt} \right)^2 - 2a \sin \theta \frac{d\omega}{dt}$$

$$\Rightarrow r \omega^2 = -2a \cos \theta \omega^2 - 2a \sin \theta \frac{d\omega}{dt}$$

$$\Rightarrow r \omega^2 = -r \omega^2 - 2a \sin \theta \frac{d\omega}{dt} \quad [r = 2a \cos \theta]$$

$$\Rightarrow 2a \sin \theta \frac{d\omega}{dt} = -2r \omega^2 \quad \Rightarrow \frac{d\omega}{dt} = -\frac{r \omega^2}{a \sin \theta}$$

$$\therefore \frac{d\omega}{dt} = \frac{-2r\omega^2}{2a \sin \theta} = \frac{-2 \cdot 2a \cos \theta \cdot \omega^2}{2a \sin \theta}$$

$$\Rightarrow \frac{d\omega}{dt} = -2\omega^2 \cot \theta. \quad \text{Proved}$$

Q A particle describes the curve $r = ae^{\theta}$ with constant angular velocity. Show that its radial acceleration is zero and the transverse acceleration varies as its distance from the pole.

Soln Given that angular velocity = constant.

$$\Rightarrow \frac{d\theta}{dt} = k \text{ (suppose)} \quad \text{--- (1)}$$

$$\because r = ae^{\theta} \quad \text{--- (2)} \Rightarrow \frac{dr}{dt} = a e^{\theta} \cdot \frac{d\theta}{dt} = ae^{\theta} \cdot k \text{ [using (1)]}$$

$$\Rightarrow \frac{dr}{dt} = ae^{\theta} \cdot k = rk \text{ [using (2)]} \quad \text{--- (3)}$$

$$\Rightarrow \frac{d^2r}{dt^2} = k \cdot \frac{dr}{dt} = k \cdot rk \text{ [using (3)]} = rk^2. \quad \text{--- (4)}$$

Now, radial acceleration = $\frac{d^2r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2$

$$= rk^2 - r \cdot k^2 \text{ [using (4) and (1)]}$$

$$= 0 \quad \text{Proved}$$

Transverse acceleration a

$$= \frac{1}{r} \frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = \frac{1}{r} \frac{d}{dt} (r^2 \cdot k) \text{ [using (1)]}$$

$$= \frac{1}{r} \frac{d}{dt} (kr^2) = \frac{1}{r} \cdot k \frac{d}{dt} (r^2) = \frac{k}{r} \cdot 2r \frac{dr}{dt}$$

$$= 2k \frac{dr}{dt} = 2k \cdot rk \text{ [from (3)]} = 2k^2r \Rightarrow \text{Tr. acc} \propto r$$